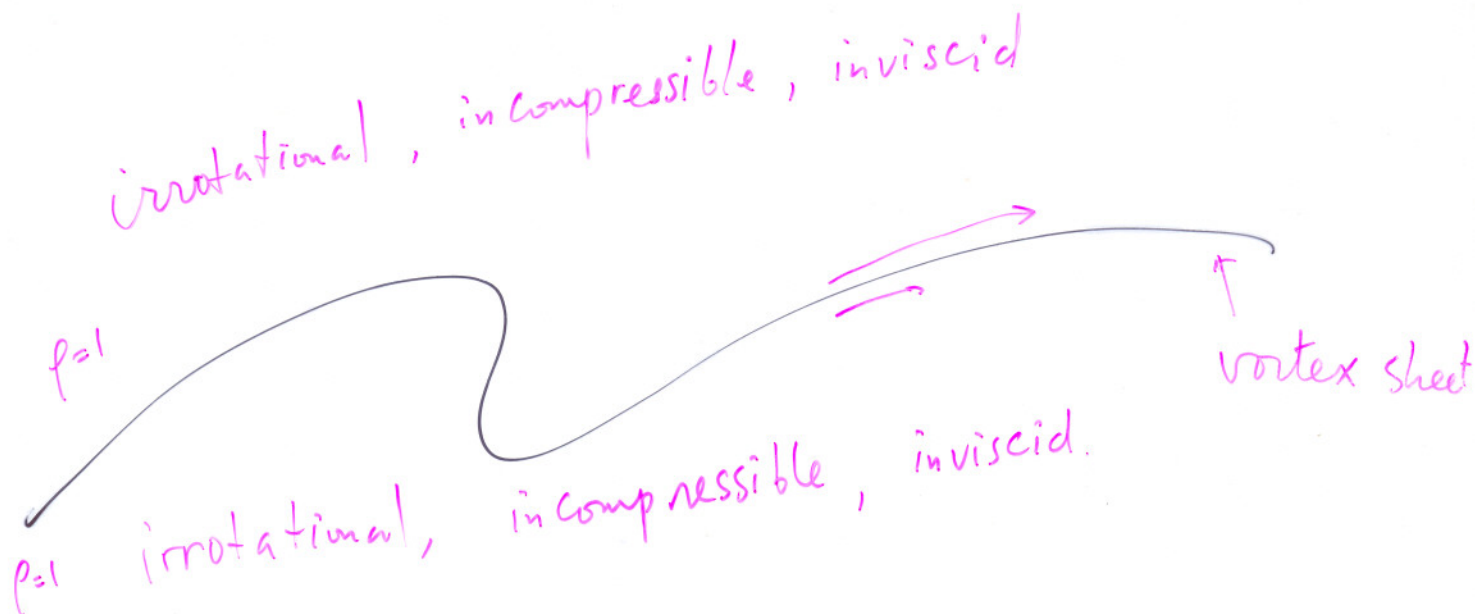


Recent Progress in Mathematical Analysis of Vortex Sheets

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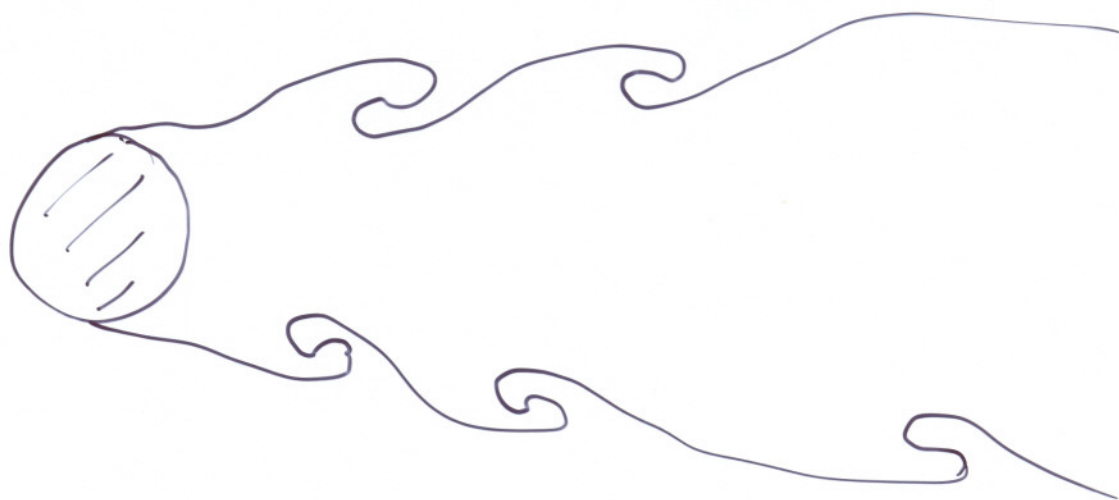
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We consider the motion of the interface separating two domains of the same fluid in \mathbb{R}^2 that moves with different velocity along the tangential direction of the interface.

this interface \rightarrow vortex sheet.

We assume the fluid occupying the two domains are incompressible, irrotational, inviscid. Also we assume the surface tension = 0.



Q: What is the specific nature of vortex sheet motion?

There are two approaches

1) Incompressible Euler Eq'n with vortex sheet initial data: let $v_0(x)$ be the initial velocity, vortex sheet data if $v_0(x)$ is discontinuous along $v_0(x)$ is a curve $\Sigma_0 \Leftrightarrow \text{curl } v_0$ is a measure supported on Σ_0 .

a). Delort (1991) Existence of weak sol'n for any initial velocity s.t. $\text{curl } v_0 \geq 0$, and $v_0 \in L^2_{loc}(\mathbb{R}^2)$

b)

b). Scheffer (1993), Shnirelman (1997, 2003)

Nonuniqueness of weak sol'n in $L^2((0, T) \times \mathbb{R}^2)$ ~~then~~

Weak sol'n provides little information about the specific nature of vortex sheet evolution.

For example: Is vortex sheet a curve at any later time if initially it is a curve?

Second approach.

2). Assuming further that vortex sheet is a curve in \mathbb{R}^2 at a later time, $t > 0$.

Incompressible Euler eq'n + this assumption

\Rightarrow Birkhoff-Rott Eq'n:

$$\bar{z}_t(\alpha, t) = \frac{1}{2\pi i} \text{P.V.} \int \frac{1}{z(\alpha, t) - z(\beta, t)} d\beta$$

(complex)

Here: $z(\alpha, t) = x(\alpha, t) + iy(\alpha, t)$ is the ^{n.t.} position of the vortex sheet at time t with the curve circulation variable α .

i.e. $\gamma = \frac{1}{|z_x|}$ is the vortex strength.

$z = z(\alpha, t) =$ (complex) position of the vortex sheet curve at time t.

Sulem, Sulem, Bardos, Frisch (1981) Existence and uniqueness of solns in analytic class for any given analytic data.

D. W. Moore (1979) analytic evidence that vortex sheet can develop singularity from analytic data.

Krasny (1986), Meiron, Baker & Orszag (1982) numerical confirmation

Duchon & Robert (1988), Caflisch & Orlandi (1989) constructed specific examples of solns of the Birkhoff-Rott eqn, where a curvature singularity develops in finite time from analytic data

\Rightarrow Birkhoff-Rott eqn is illposed in $C^{1+p}, H^s(\mathbb{R})$
 $p > 0, s > 3/2.$

G. Lebeau (2001)

Q. Are there solns in a class larger than C^{1+p} or H^s , ($p > 0, s > 3/2$)?

What is the nature of the vortex sheet at and beyond the singularity formation time?

- We look for solns of the B-R Eqn in the largest possible spaces where the eqn makes sense.
- Experimental & numerical evidences suggest that this class should include ^(infinite) spirals.
- The class C^{1+p} or H^s , $p > 0$, $s > 3/2$ does not contain ^(infinite) spiral curves.

These lead us to consider solns in Chord-arc class.

- Chord-arc class contains log spirals.
- Chord-arc class is almost the largest in L^2 sense.

Chord-arc curves : (non self intersecting)

Let $\xi = \xi(s)$ describes ~~the~~ ^a curve Γ in arclength s .



We say Γ is chord-arc, if $\exists M > 1$. s.t.

$$|s_1 - s_2| \leq M |\xi(s_1) - \xi(s_2)| \quad \text{for all } s_1, s_2.$$

\uparrow
arclength

\uparrow
chord-length

The infimum of all such M is called the chord-arc constant.

$$\text{Def: } m_Q(f) = \frac{1}{|Q|} \int_Q f(x) dx$$

Let $I \subset \mathbb{R}$ be an interval, We say a fcn $f \in L^1_{loc}(\mathbb{R})$ is in $BMO(I)$ if

$$\|f\|_{BMO(I)} = \sup_{Q \subset I, Q \text{ interval}} \frac{1}{|Q|} \int_Q |f(x) - m_Q f| dx < \infty.$$

(G. David)

Thm: 1) If $b \in BMO(\mathbb{R})$, and $\|b\|_{BMO} < 1$,

then $\xi(s) = \xi_0 + \int_0^s e^{ib(s')} ds'$ defines a chord-arc curve, with chord-arc constant $\leq \frac{1}{1 - \|b\|_{BMO}}$.

2) If $\xi = \xi(s)$ defines a chord-arc curve, s -arclength,

then $\xi'(s)$ exist a.e. and $\exists b \in BMO$ s.t.

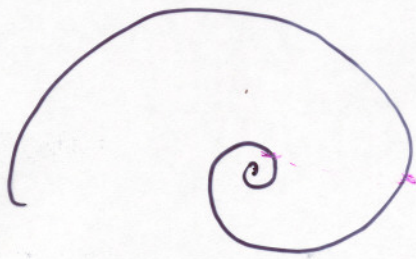
$$\xi'(s) = e^{ib(s)} \quad \text{Moreover, if the chord-arc constant}$$

$$M \approx 1, \text{ then } \|b\|_{BMO} \approx 0.$$

Example: take $b = \varepsilon \ln|s| \in BMO$, $\varepsilon \ll 1$.

$$\text{Let } \xi(s) = \int_0^s e^{i\varepsilon \ln|s'|} ds' = \int_0^s |s'|^{i\varepsilon} ds' = \frac{1}{1+i\varepsilon} s^{1+i\varepsilon} \quad (s>0)$$

\neq this is a log spiral.



Thm: The Birkhoff-Rott Eqn is of "elliptic type" in the chord-arc class.

A typical elliptic eqn: $\Delta u = u_{tt} + u_{xx} = 0$

Initial value problem: (*) $\begin{cases} u_{tt} + u_{xx} = 0 \\ u(x, 0) = u_0(x) \\ u_t(x, 0) = u_1(x) \end{cases} > \text{given}$

$\hat{u}(\xi, t)$ = Fourier transform of $u(x, t)$ in variable x

Soln $\hat{u}(\xi, t) = \frac{1}{2} (\hat{u}_0(\xi) + \frac{1}{|\xi|} \hat{u}_1(\xi)) e^{|\xi|t} + \frac{1}{2} (\hat{u}_0(\xi) - \frac{1}{|\xi|} \hat{u}_1(\xi)) e^{-|\xi|t}$

\Rightarrow Conclusions:

1) (*) can be solved in $L^2(\mathbb{R})$ for $0 \leq t \leq T$, if and only if $\hat{u}_0(\xi) + \frac{1}{|\xi|} \hat{u}_1(\xi)$ is real analytic

i.e. $\exists a > 0$ s.t. $(\hat{u}_0(\xi) + \frac{1}{|\xi|} \hat{u}_1(\xi)) e^{a|\xi|} \in L^2(\mathbb{R})$

2) If for $0 \leq t \leq T$, we have $u(\cdot, t) \in L^2(\mathbb{R})$

then $(\hat{u}_0(\xi) + \frac{1}{|\xi|} \hat{u}_1(\xi)) e^{T|\xi|} \in L^2(\mathbb{R}) \Rightarrow$

$(\hat{u}_0(\xi) + \frac{1}{|\xi|} \hat{u}_1(\xi)) e^{t|\xi|} = (\hat{u}_0(\xi) + \frac{1}{|\xi|} \hat{u}_1(\xi)) e^{T|\xi|} e^{(t-T)|\xi|}$ is analytic for $0 < t < T$.

$\Rightarrow u(\cdot, t)$ real analytic for $0 < t < T$.

Thm 1: Assume that $Z \in H^1(0, T], L^2_{loc}(\mathbb{R}) \cap L^2(0, T; H^1_{loc}(\mathbb{R}))$

is a soln of the Birkhoff-Rott eqn for $0 \leq t \leq T$, satisfying that

1. $\exists m, M > 0$, indep of t , s.t.

$$m|\alpha - \beta| \leq |Z(\alpha, t) - Z(\beta, t)| \leq M|\alpha - \beta| \quad \forall \alpha, \beta \in \mathbb{R} \\ 0 \leq t \leq T$$

\Rightarrow Then there is a constant $C(m, M) > 0$ as follows.

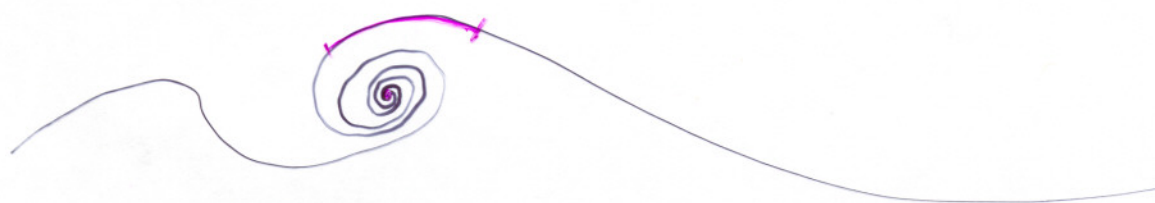
If also

2. on some fixed interval (a, b) , $\exists \delta_0 > 0$, satisfying

$$\sup_{[0, T]} \| \ln Z_\alpha(\cdot, t) \|_{BMO(a, b), \delta_0} \leq C(m, M)$$

Then in fact $Z_\alpha \in C((a, b) \times (0, T))$, and for each $t_0 \in (0, T)$

$Z_\alpha(\cdot, t_0)$ is analytic on (a, b) .



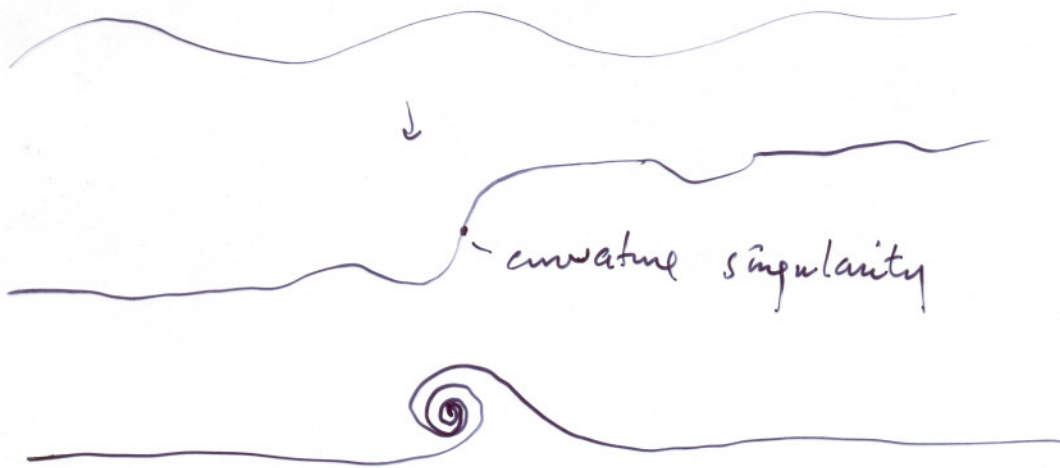
1. means: $Z = Z(\alpha, t)$ is chord-arc and the

$$\text{vortex strength } 0 < \frac{1}{M} \leq \gamma = \frac{1}{|Z_\alpha|} \leq \frac{1}{m} < \infty$$

2. \Rightarrow The vortex sheet curve does not roll-up too fast on a section (a, b) .

G. Lebeau (2001). Assuming $Z \in C^{1+p}$, $p > 0$, (stronger assumption)

Cor.: After the singularity formation time, the vortex sheet will not be a chord-arc curve that doesn't roll-up too fast.



Initial value problem for the Birkhoff-Rott eq'n.

If the initial velocity $v_0(x)$ is given, then the position of the initial vortex sheet is given, (where the discontinuity is.)

and the initial vortex strength γ is also given.

$\Rightarrow z_0 = z(\alpha, 0)$ is completely given.

$\Rightarrow \ln z_\alpha(\alpha, 0)$ is given.

(+i) $\ln z_\alpha(\alpha, 0) = \text{Real part} + i \text{Imaginary part}$

Thm 2: For any real valued fcn $w_0 \in H^{3/2}(\mathbb{R})$,
 there exist a $T = T(\|w_0\|_{H^{3/2}}) > 0$, such that the
 Birkhoff-Rott eq'n has a sol'n $z = z(\alpha, t)$ for $0 \leq t \leq T$,
 satisfying $\text{Im} \{ (1+i) \ln z_\alpha(\alpha, 0) \} = w_0(\alpha)$, and
 $z(\cdot, t)$ is real analytic for each $0 < t < T$.

Thm 3: Assume that $z \in H^1([0, T], L^2_{loc}(\mathbb{R}) \cap L^2(0, T, H^1_{loc}(\mathbb{R}))$
 is a sol'n of the Birkhoff-Rott eq'n for $0 \leq t \leq T$,
 $T > 0$, satisfying properties 1 and 2 of Thm 1.
 on some interval (a, b) . Assume further that
 $w_0 = \text{Im} \{ (1+i) \ln z_\alpha(\cdot, 0) \}$ is analytic on (a, b) ,
 Then $z_\alpha \in C((a, b) \times [0, T])$, and $\text{Re} \{ (1+i) \ln z_\alpha(\cdot, 0) \}$
 is also analytic on (a, b) .